

**LANL/IGPP Climate Study Group**  
**04/10/2007**

# **Three-Dimensional Radiative Transfer:**

*What is it?*

**&**

*Why does it matter?*

**Anthony B. Davis**

*Los Alamos National Laboratory*

**Space & Remote Sensing Group (ISR-2)**



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**LA-UR 07-2467**



# **... with ideas, papers, discussions, critiques from/by many colleagues. For starts ...**

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Qi-Long Min (SUNY-Albany)

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Greame Stephens (CSU)

Dennis O'Brien (CSU)

Alex Kostinski (MTU)

Mike Larson (ARL)

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Tamas Várnai (NASA-GSFC)

Lazaros Oreopoulos (GSFC)

Frank Evans (U of Co)

Bernard Mayer (DRL)

Eugene Clothiaux (PSU)



## Over the past 9+ years, funding from:

- \* DOE's Multispectral Thermal Imager (MTI)
- \* NNSA/NA-22 programs
- \* ISR Division's PD- and TD-programs
- \* LDRD/ER, -/DR, and -/PRD
- \* NASA-GSFC Collaborations ("work of others")
- \* [SC/BER] ARM Program - Science Team PI



# Prologue

**... followed by 4 acts  
&  
an epilogue**

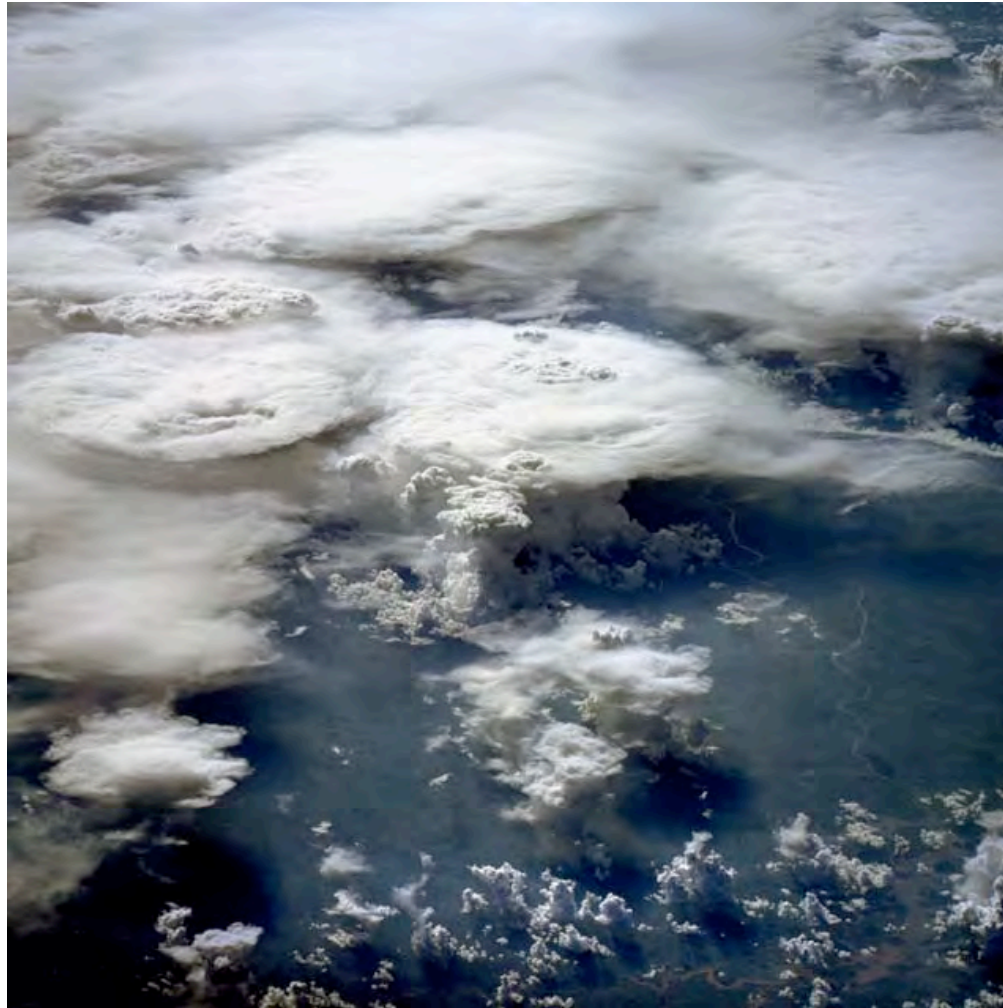
# State-of-the-Art Conceptual Model for Radiative Transfer in Clouds

- To compute radiative fluxes inside most Global Climate Models (GCMs) radiation modules, as well as cloud system resolving models (CSRMs) and even cloud process models.
- To compute radiances, hence physical cloud properties, in *operational* cloud remote sensing schemes at NASA and Co, irrespective of pixel size.



*This is a cloud!*

# Reality:



# What we\* compute versus what we\* measure

$I_\lambda(\mathbf{x}, \vec{\Omega})$ , from 3D Radiative Transfer Equation (RTE), plus BCs.

$$\Phi = \int_{\lambda_{\min}}^{\lambda_{\max}} f(\lambda) \int_{\mathbf{x} \in S \subseteq \mathcal{R}^3} \chi(\mathbf{x}) \int_{4\pi} r(\vec{\Omega}) I_\lambda(\mathbf{x}, \vec{\Omega}) d\lambda d\mathbf{x} d\vec{\Omega}$$

E.g., BB actinic flux in a CSRM or LES cell :

$$\lambda_{\min} = 0, \lambda_{\max} = \infty, f(\lambda) \equiv 1$$

$$\chi_{\mathbf{x}^*}(\mathbf{x}) = \begin{cases} 1 / \Delta x \Delta y \Delta z, & x^* \leq x < x^* + \Delta x, y^* \leq y < y^* + \Delta y, z^* \leq z < z^* + \Delta z \\ 0, & \text{otherwise} \end{cases}$$

$$r(\vec{\Omega}) \equiv 1$$

**\* “we” can be Nature**

# 3D RT Problem Space

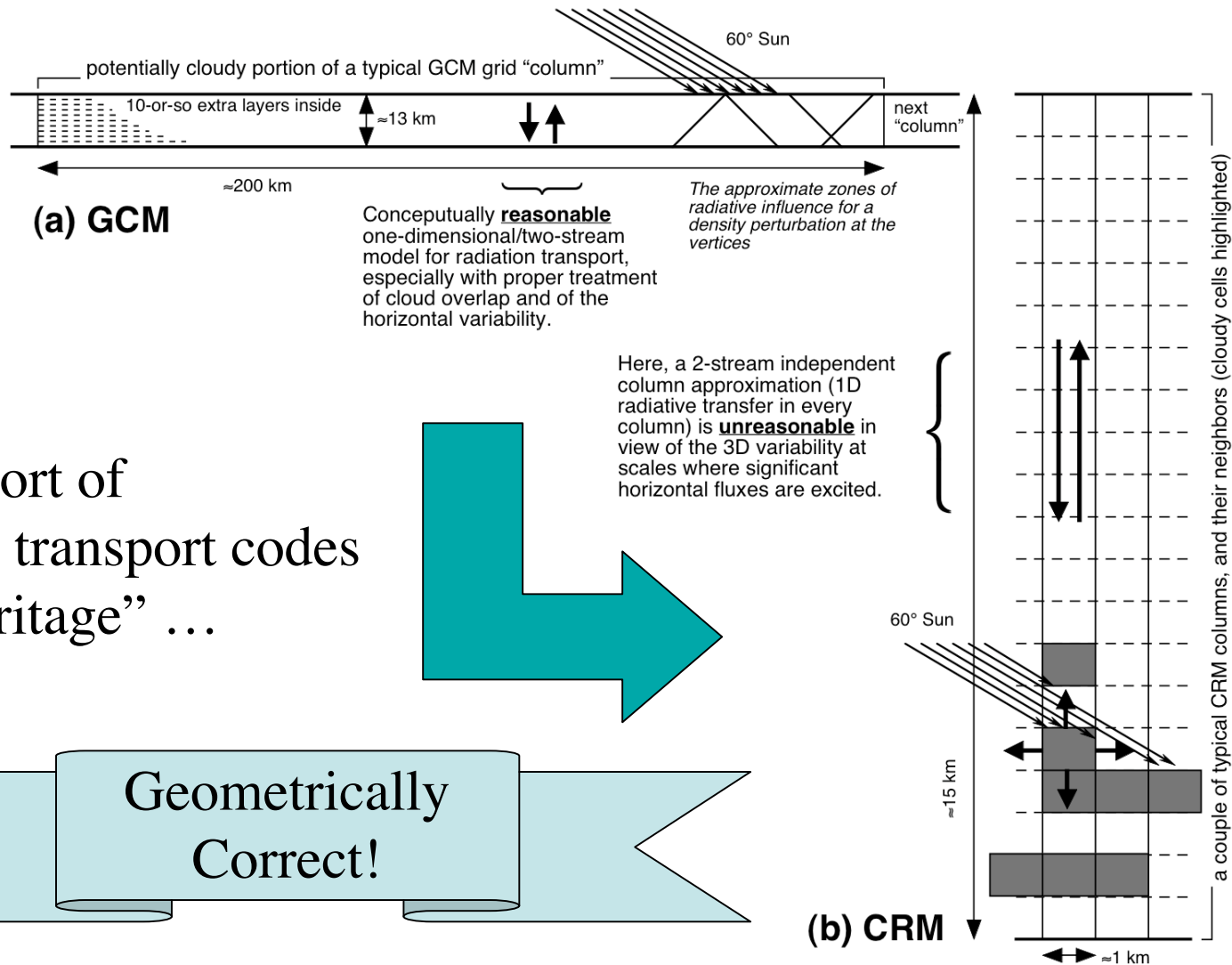
$I_\lambda(\mathbf{x}, \vec{\Omega})$ $f(\lambda)\chi(\mathbf{x})r(\vec{\Omega})$	<b>Diagnostics</b> using radiances [samples of $\Omega$ ]	<b>Energetics</b> using fluxes [sums over $\Omega$ ]
“pixel” scales: <u>resolved</u> structure [samples of $\mathbf{x}$ ]	“Pixel/column adjacency” problems <b>I3RC</b>	3D radiative heating/cooling rates in CSRMs <b>I3RC-Approx</b>
“domain” scales: <u>unresolved</u> structure [sums over $\mathbf{x}$ ]	Large-footprint (or beam-filling) problems <b>RAMI</b>	GCM radiation parameterization schemes <b>ICRCCM</b>



# ACT 1

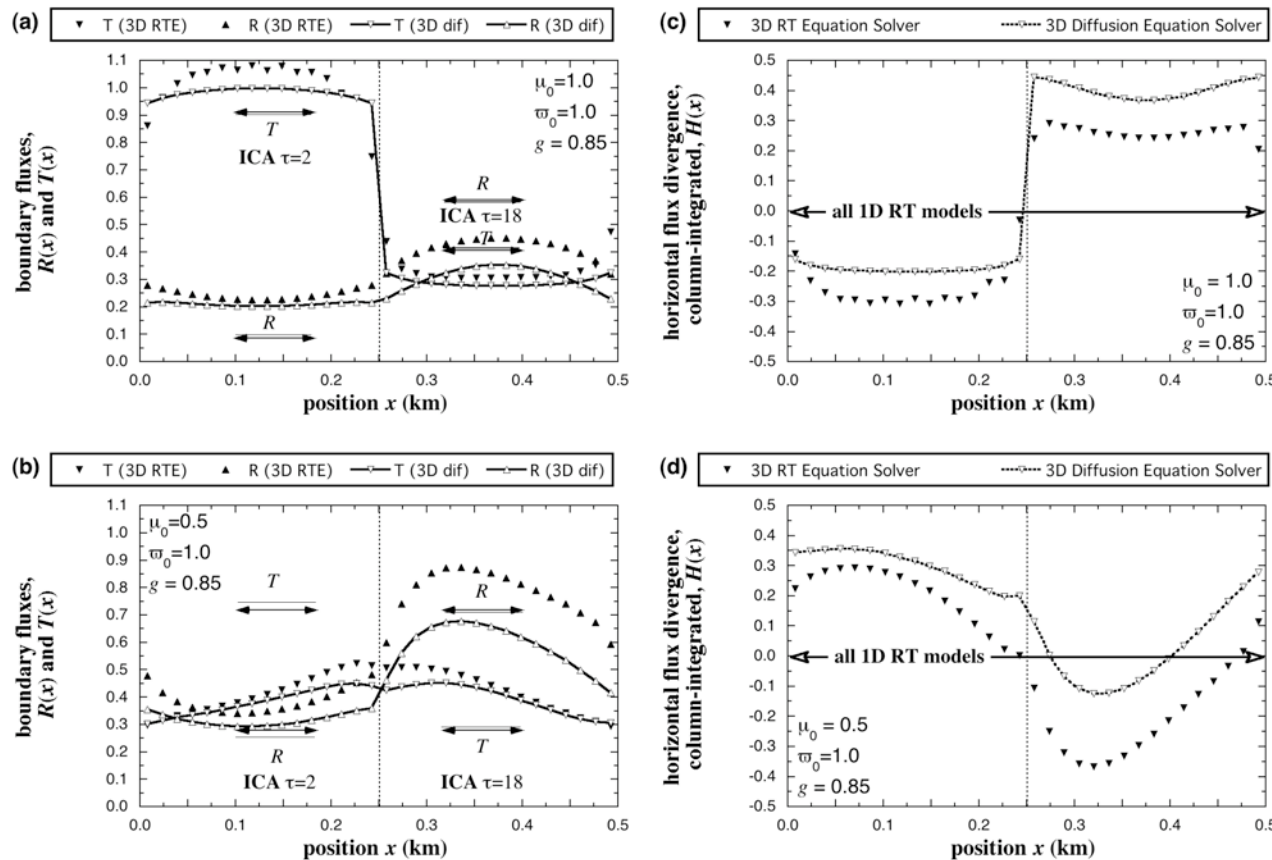
**“Good enough” 3D RT**

# In 3D RT, consider the inner and outer aspect ratios.



# 1D Independent Column Approximation (ICA) versus 3D Transport and Diffusion

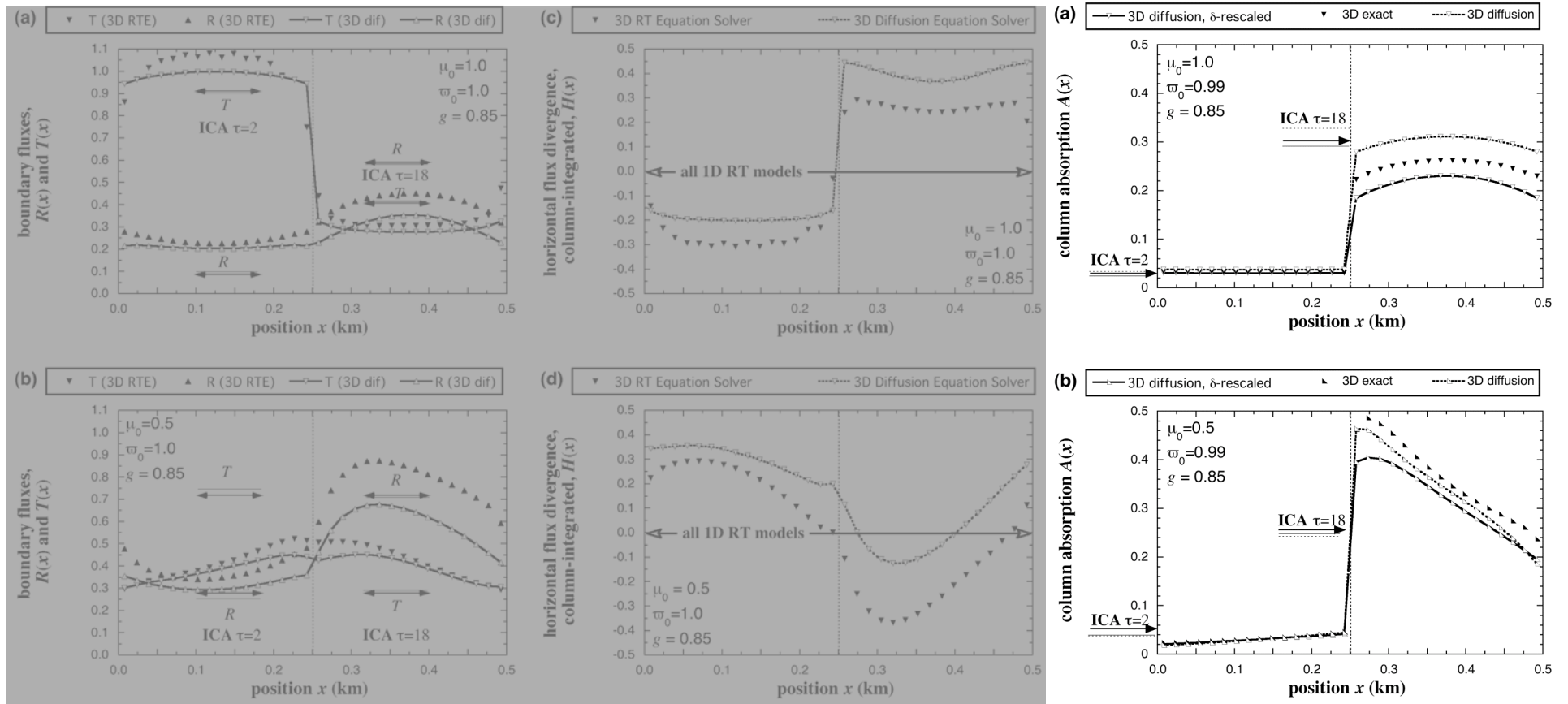
I3RC “Case 1” square-wave cloud: 0.25 km thick, alternating  $\tau = 2, 18$  every 0.25 km.



$$\varpi_0 = 1$$

# 1D Independent Column Approximation (ICA) versus 3D Transport and Diffusion

I3RC “Case 1” square-wave cloud: 0.25 km thick, alternating  $\tau = 2, 18$  every 0.25 km.



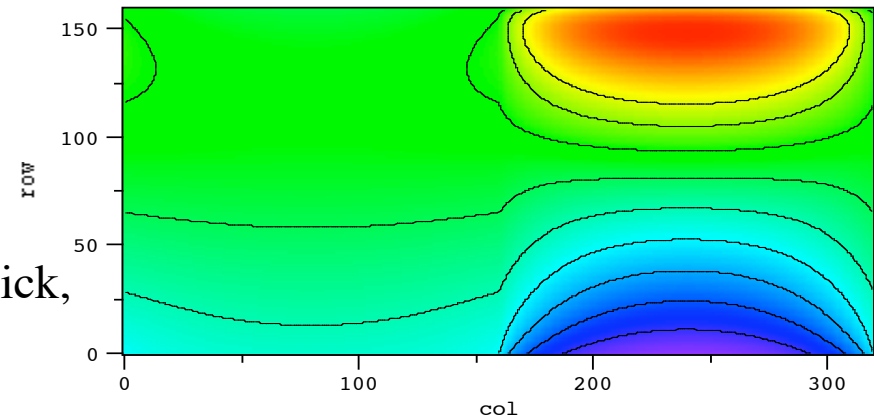
$\tau_0 = 1$

$\tau_0 = 0.99$

# 3D RT in $O(N)$ FLOPs: Computational Diffusion Theory

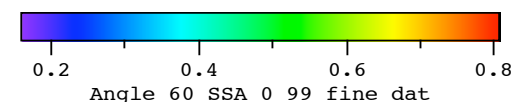
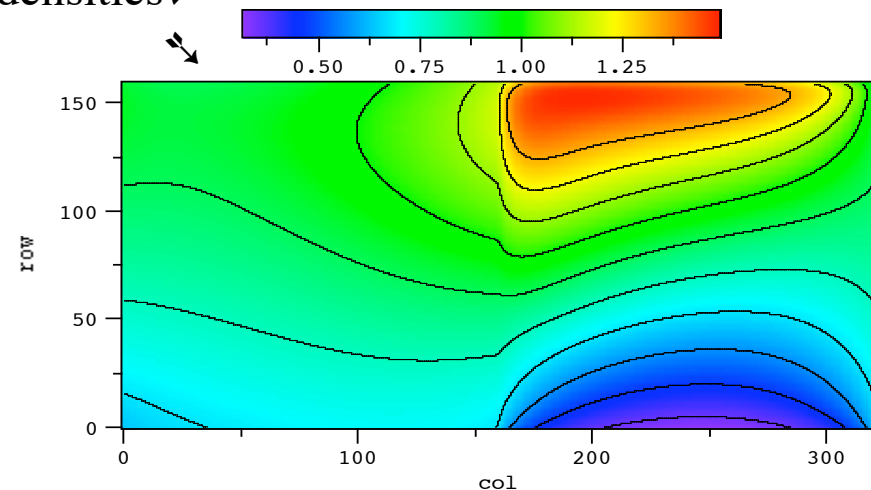
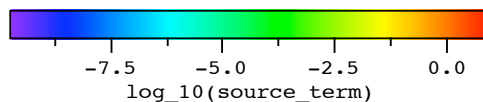
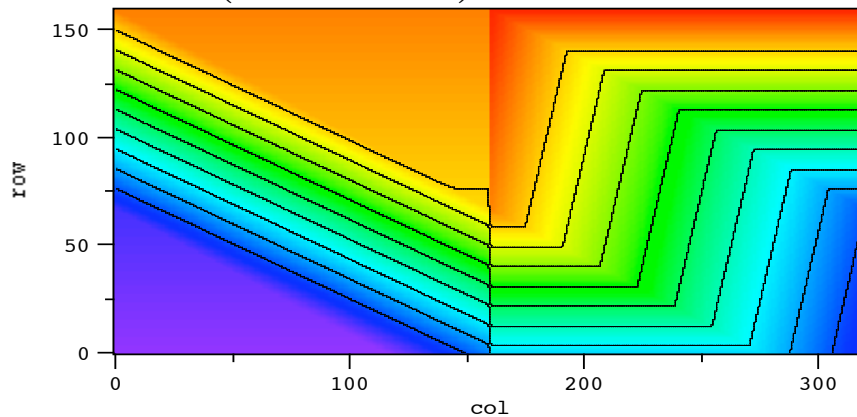
- Joint work with Mike Hall (CCS-2)
- Ideal for stratus layers
- Multi-grid solver, parallelizing code

I3RC “Case 1” square-wave cloud: 0.25 km thick,  
alternating  $\tau = 2, 18$  every 0.25 km.



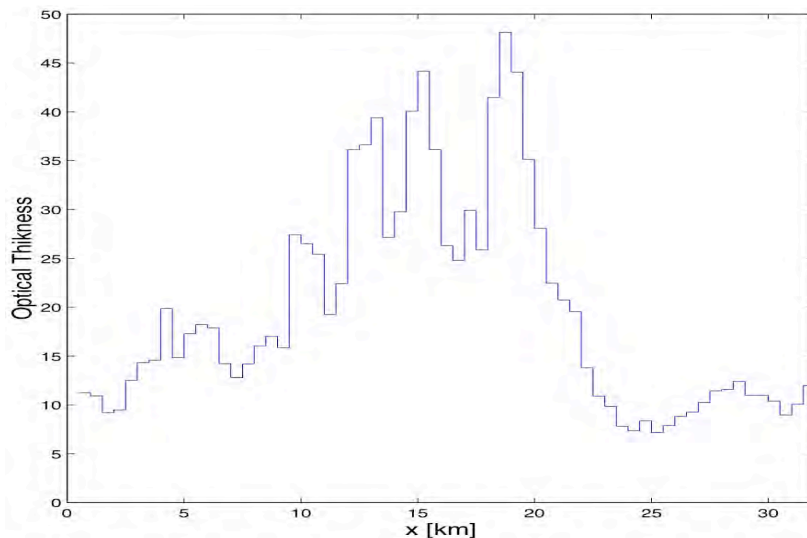
Diffuse photon densities ↗

Direct flux (source term)

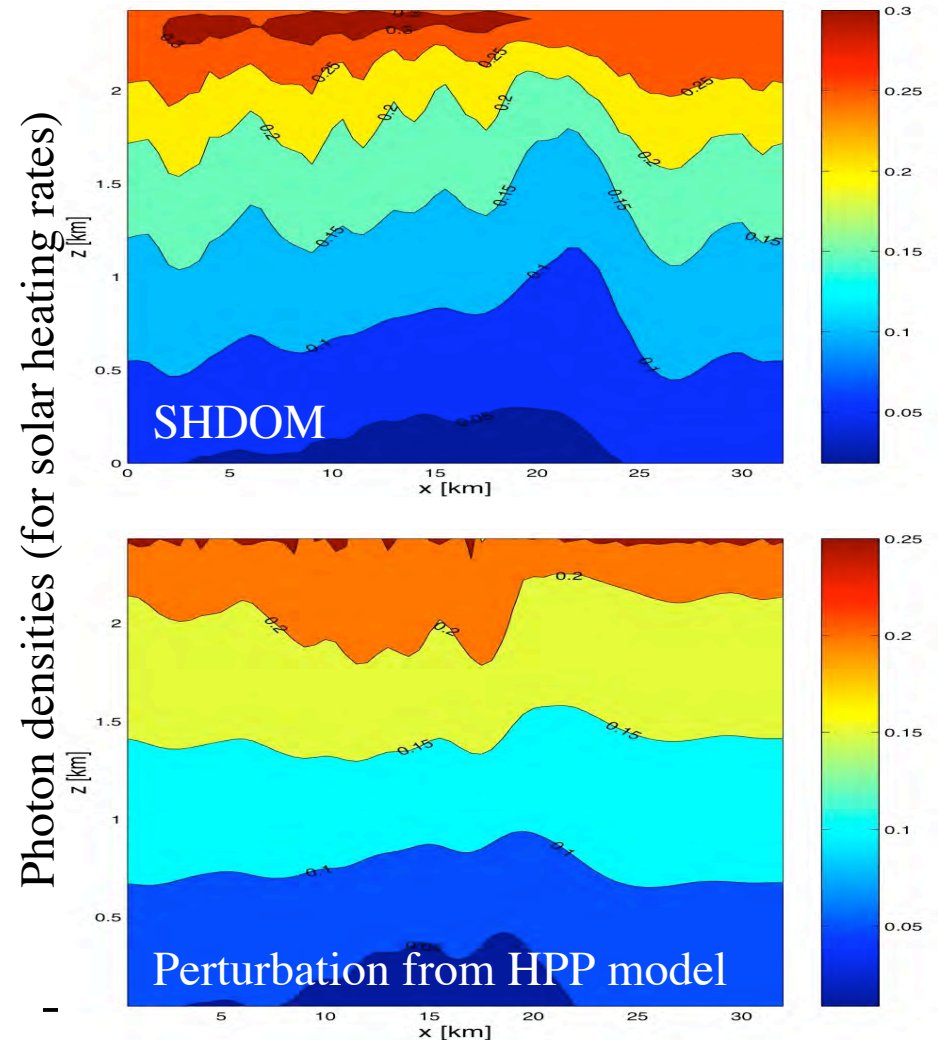


# 3D RT in Closed-Form Expressions: Adjoint Perturbation Theory

I3RC “Case 2” Cloud from MMCR+



- Joint work with Igor Polonsky (now at CSU)
- Uses spatial Green functions for Homogeneous Plane-Parallel (HPP) clouds
- Perturbation could depart from
  - ICA, rather than HPP model (done)
  - computational 3D diffusion model (cf. previous slide)



# ACT 2

**Passive from above,  
active from below**

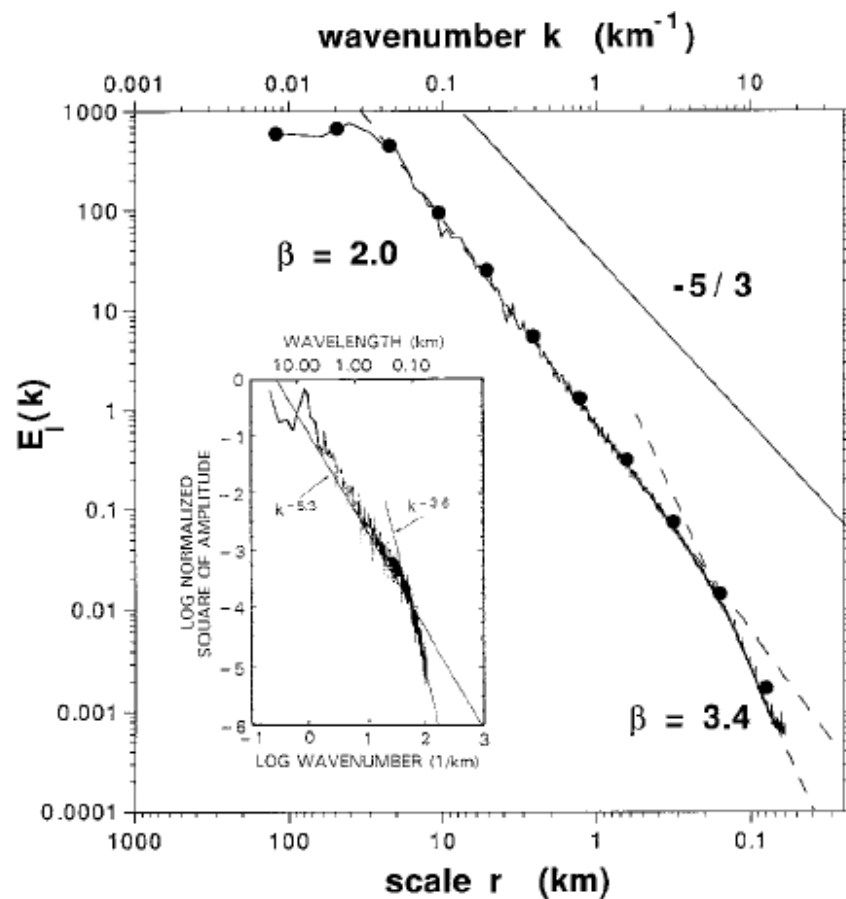
# Context, c. 1995

## The Landsat Scale Break in Stratocumulus as a Three-Dimensional Radiative Transfer Effect: Implications for Cloud Remote Sensing

ANTHONY DAVIS,\* ALEXANDER MARSHAK,\* ROBERT CAHALAN, AND WARREN WISCOMBE

*Climate and Radiation Branch, NASA/Goddard Space Flight Center, Greenbelt, Maryland*

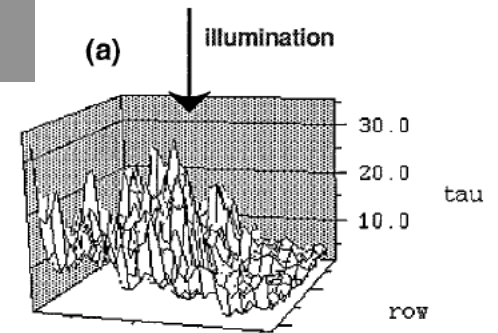
(Manuscript received 28 August 1995, in final form 16 May 1996)



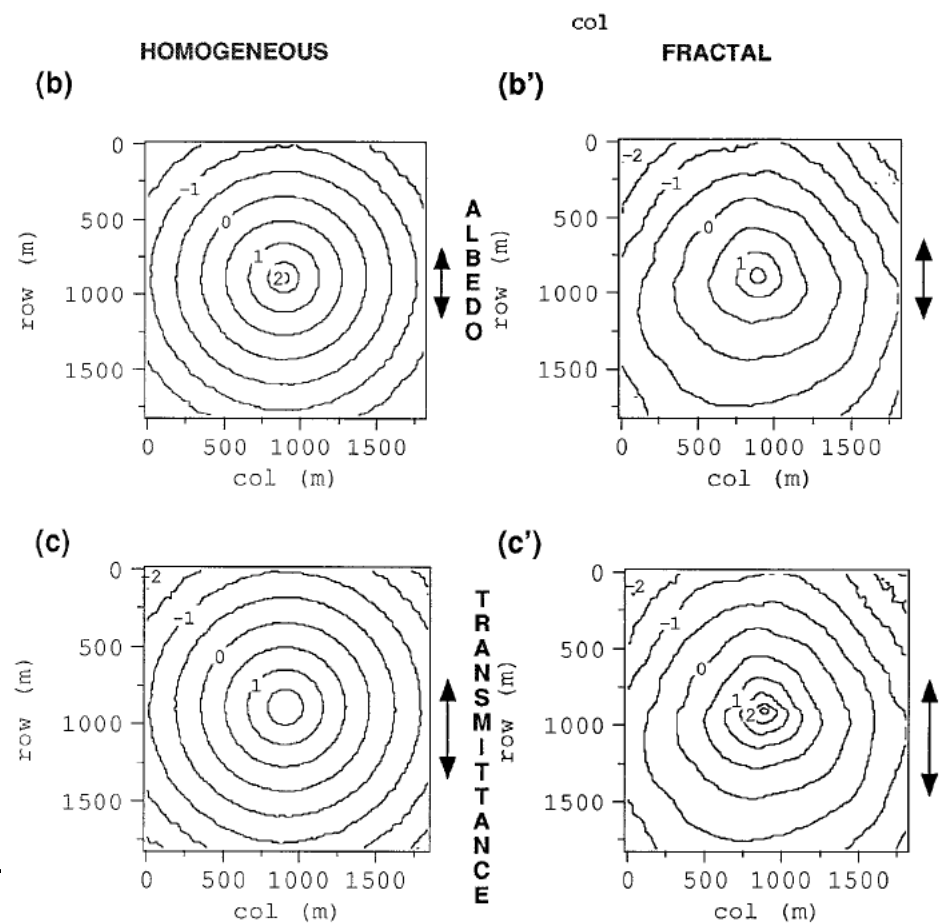
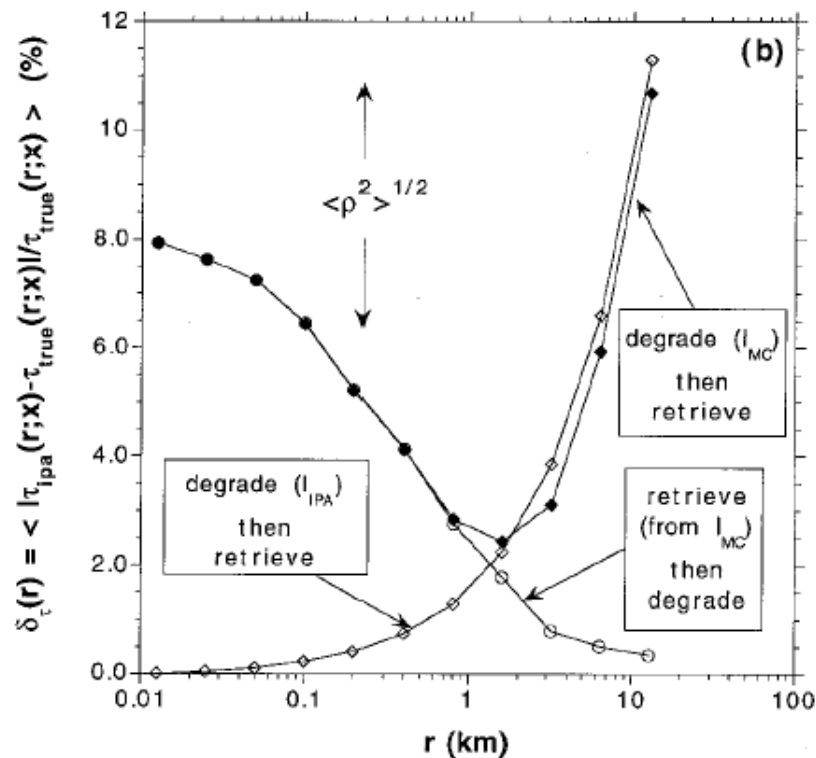


# Context, c. 1995

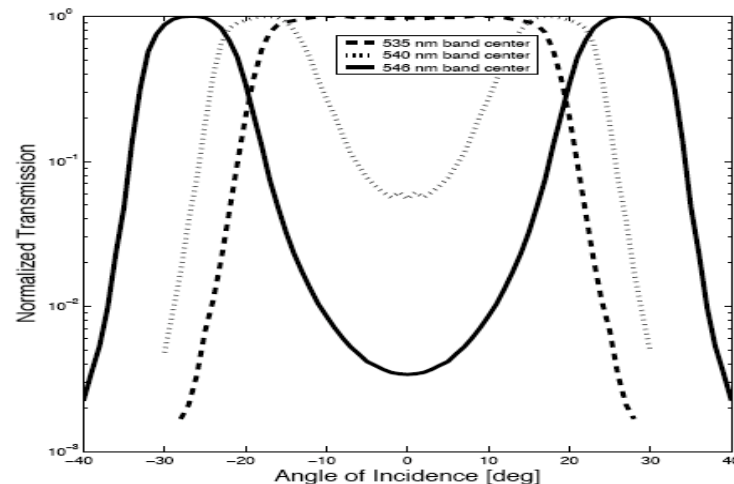
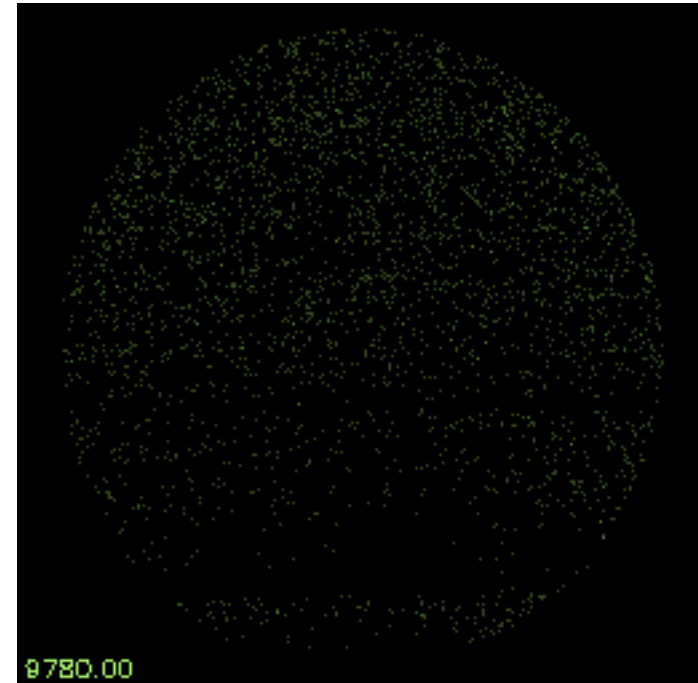
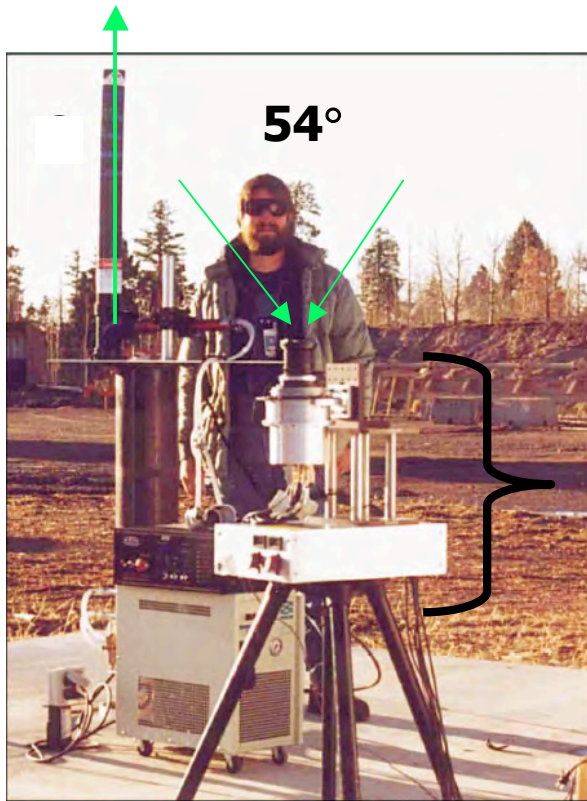
“Sasha, trust me!”



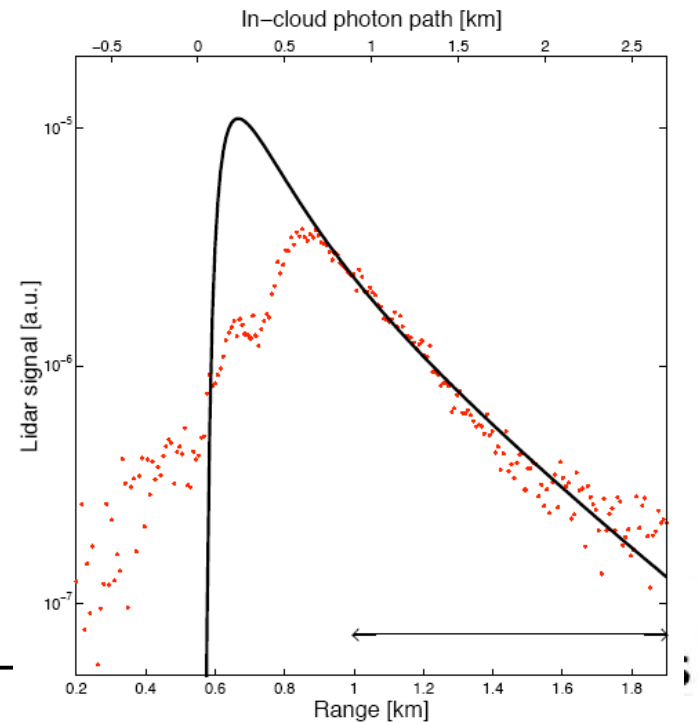
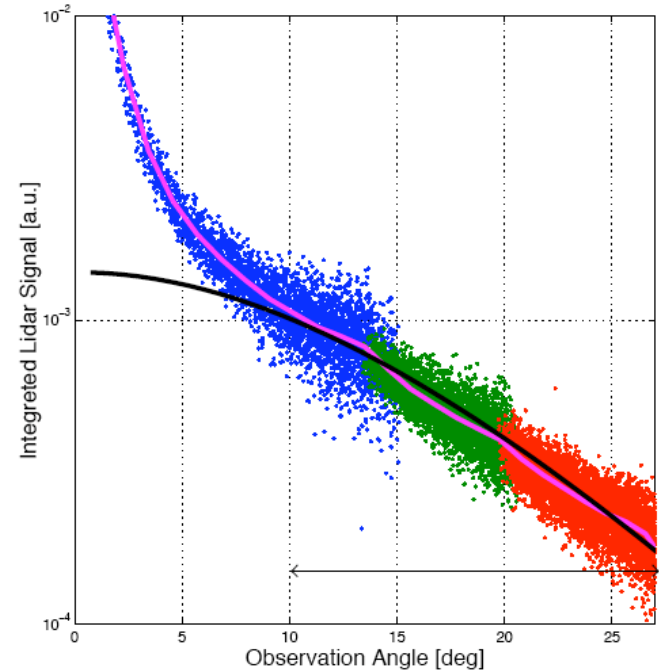
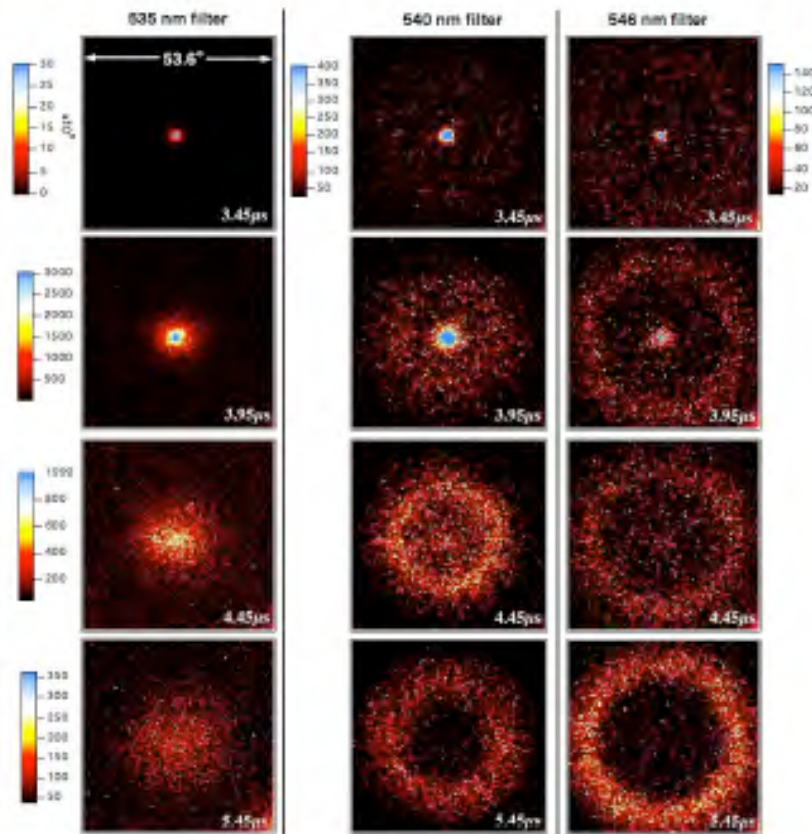
“No (pixel) scale left behind!”



# Wide-Angle Imaging Lidar (WAIL)

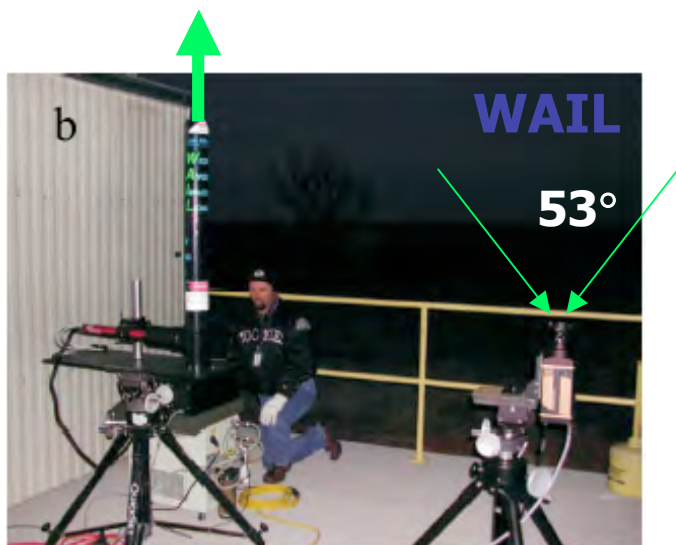


# Wide-Angle Imaging Lidar (WAIL-2): New detector (gated/intensified CCD)

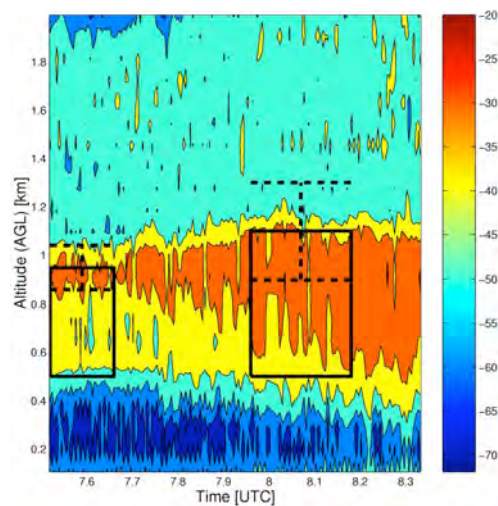


# Wide-Angle Imaging Lidar (WAIL) validation campaign at ARM SGP site

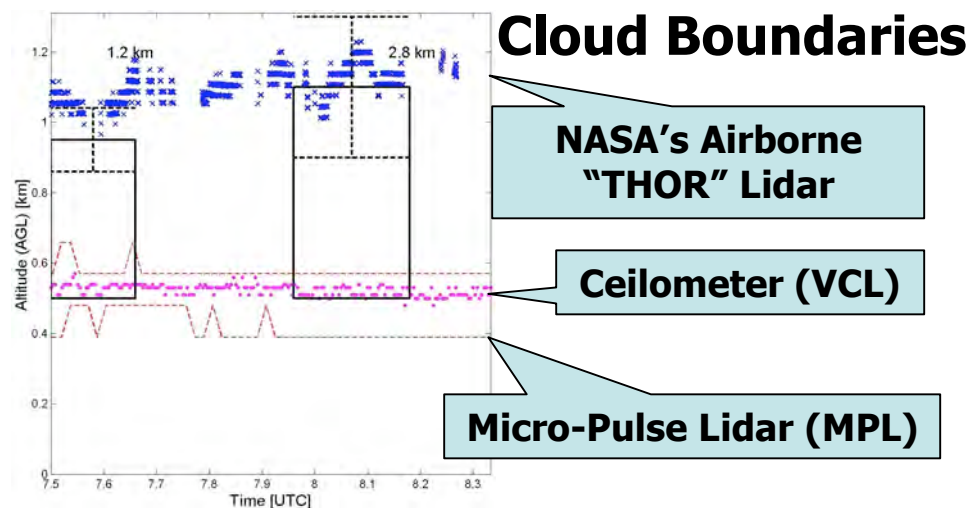
Polonsky,  
Davis, Love



... captures and analyses the space-time  
Green function, in green light!



**mm-Radar  
Reflectivity  
(MMCR)**

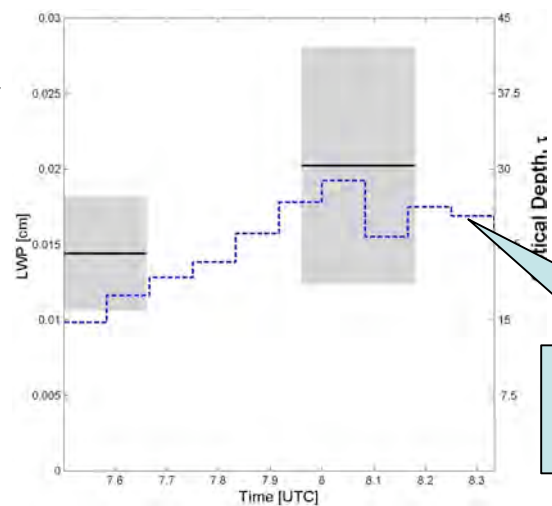


**Cloud Boundaries**

NASA's Airborne  
"THOR" Lidar

Ceilometer (VCL)

Micro-Pulse Lidar (MPL)

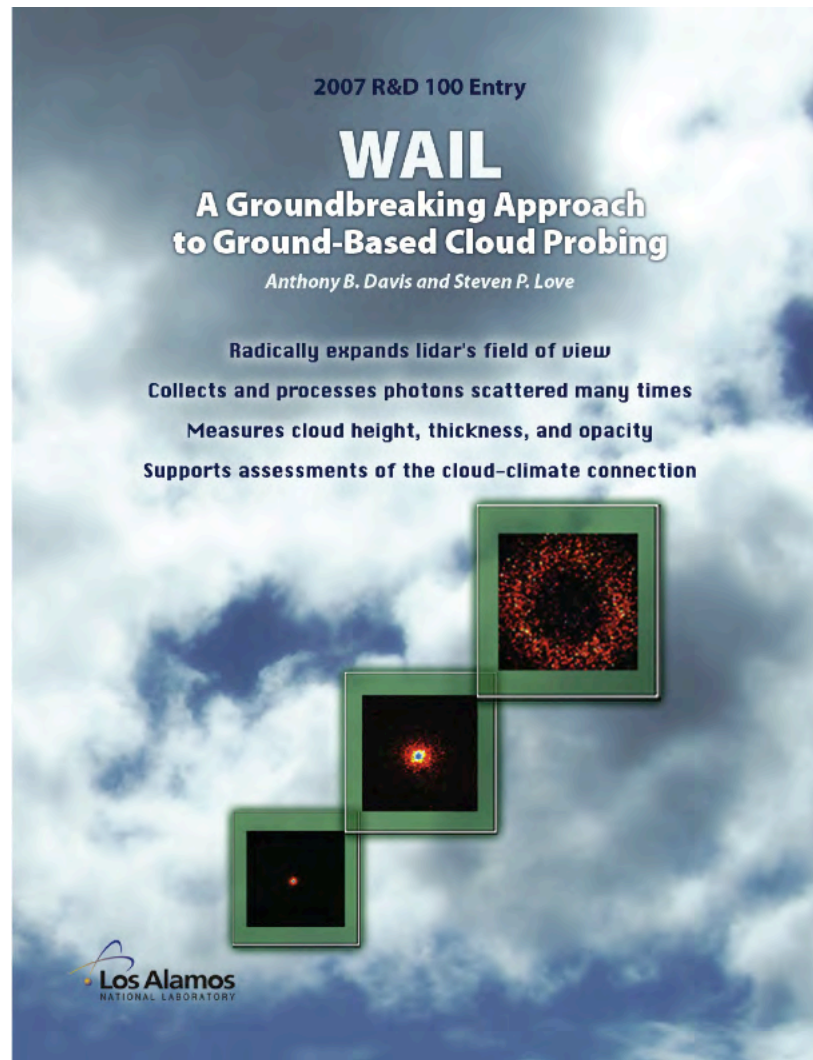


**Cloud Optical  
Depth**

2-channel  $\mu$ wave  
radiometer (MWR)



# Wish us luck ...

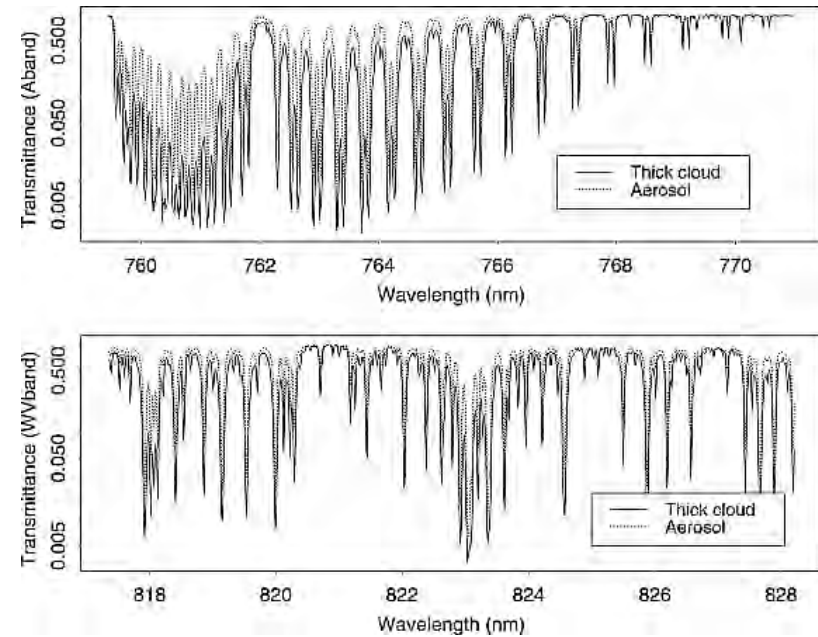


# ACT 3

**RT in 3+1 dimensions,  
using a steady source**

# Differential absorption spectroscopy at high resolution

From: Min Q.-L., L. C. Harrison, P. Kiedron, J. Berndt, and E. Joseph, 2004: A high-resolution oxygen A-band and water vapor band spectrometer, *J. Geophys. Res.*, **109**, D02202, doi:10.1029/2003JD003540.



	x-section	density	pathlength	
	$\sigma_v$	$n$	$L$	$I(v)/I_0 = \exp[-\sigma_v \times n \times L]$
known/not :	?	✓	✓	estimating molecular cross - sections in the laboratory
	✓	?	✓	monitoring amounts of chemical effluent in situ
	✓	✓	?	scattering/reflection diagnostics of media permeated with gas

$$I(v) = I(k_v) = I_0 \int_0^{\infty} \exp[-k_v L] p(L) dL \text{ (equivalence "theorem")} \Rightarrow \langle L \rangle = - \left( \frac{d}{dk_v} \right) \ln I(k_v)$$

Temporal Green function of “3+1 D” RTE

# Time-dependent diffusion theory for transmitted fluxes: Plane-parallel cloud model

$H$ : cloud thickness

$\tau$ : cloud optical thickness ( $\sigma H = H / \text{mean-free-path}$ )

$g$ : asymmetry factor of scattering phase function

(i.e.,  $\langle \cos \theta_s \rangle = 2\pi \int \cos \theta_s p(\cos \theta_s) \sin \theta_s d\theta_s = 0.75 - 0.85$ )

**From asymptotic theory:**

*(i.e., scaling arguments based on random walk statistics)*

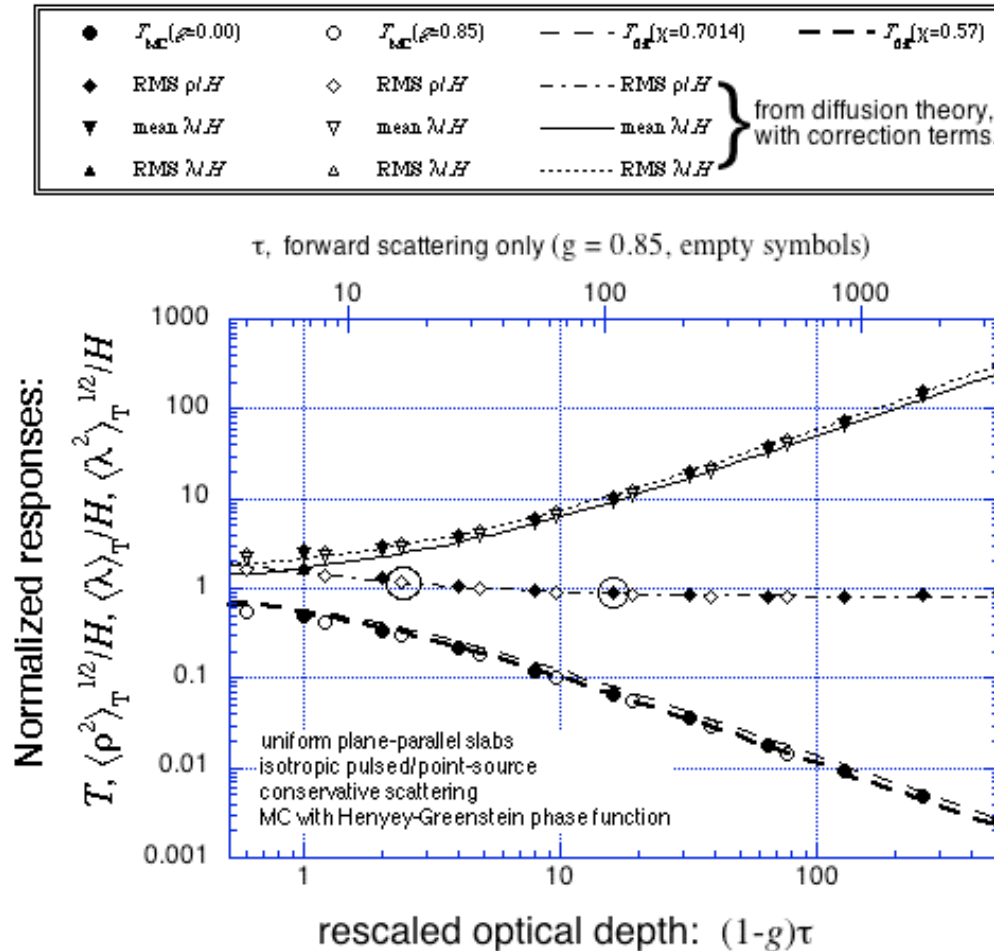
$$\langle L \rangle_T \propto (1 - g)\tau \times H$$

$$\langle \rho^2 \rangle_T^{1/2} \propto H^2$$

$$T \propto \frac{1}{(1 - g)\tau}$$



# Exact diffusion theory vs. Monte Carlo?

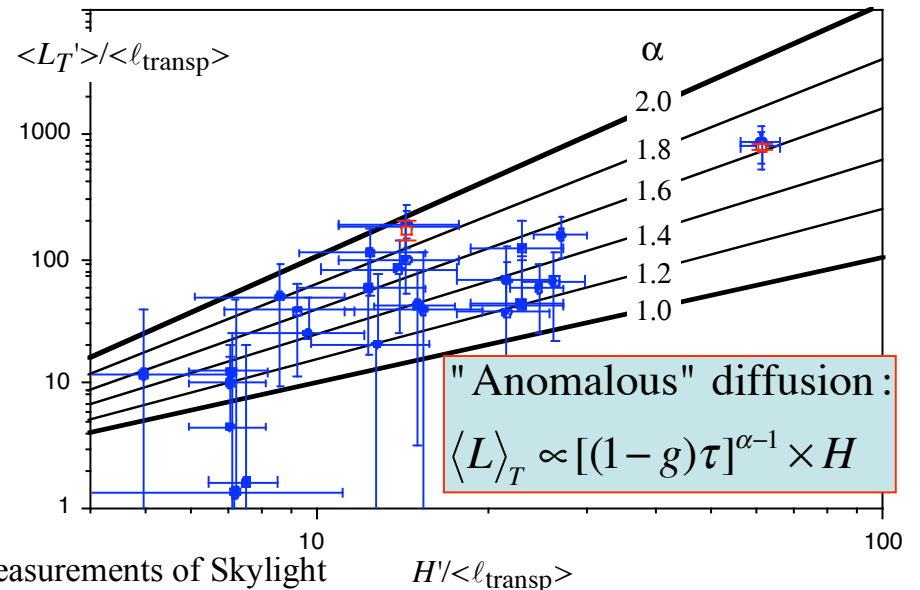


$$\langle L^2 \rangle_T^{1/2} / \langle L \rangle_T \approx \sqrt{7/5} \approx 1.18$$

Davis, A. B., and A. Marshak, 2002: Space-time characteristics of light transmitted by dense clouds, A Green function analysis, *J. Atmos. Sci.*, **59**, 2713-2727.

# Ground-based oxygen spectroscopy

Cases near the  $\alpha=2$  line are very overcast, and those near  $\alpha=1$  are for sparse clouds, as expected from model.

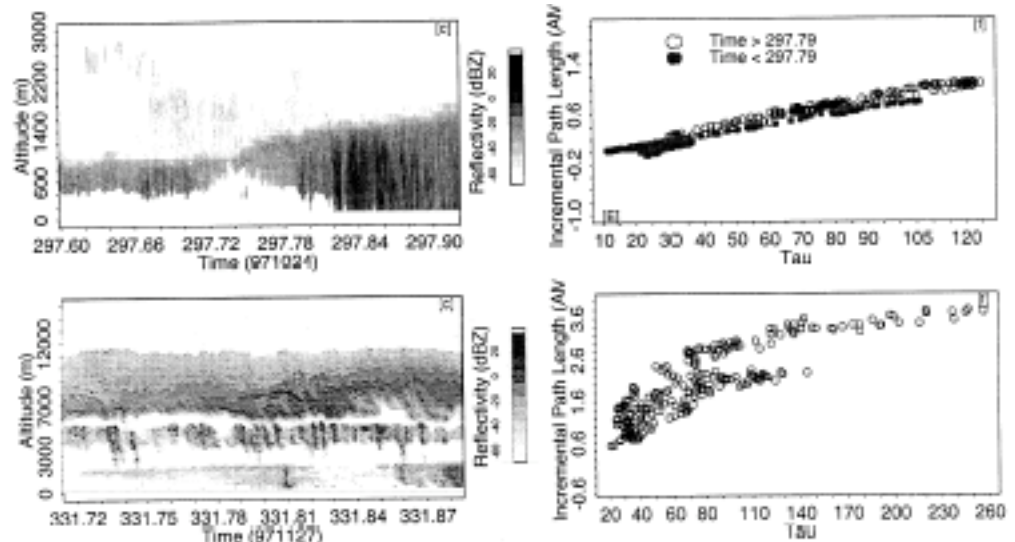


Pfeilsticker, K., 1999: First Geometrical Pathlength Distribution Measurements of Skylight Using the Oxygen A-band Absorption Technique - II, Derivation of the Lévy-Index for Skylight Transmitted by Mid-Latitude Clouds, *J. Geophys. Res.*, **104**, 4101-4116.

(Davis & Marshak, 1997)

A single cloud layer ( $\alpha=2$ ) with variable thickness  $H \propto$  the slope of the linear path vs optical depth plot.

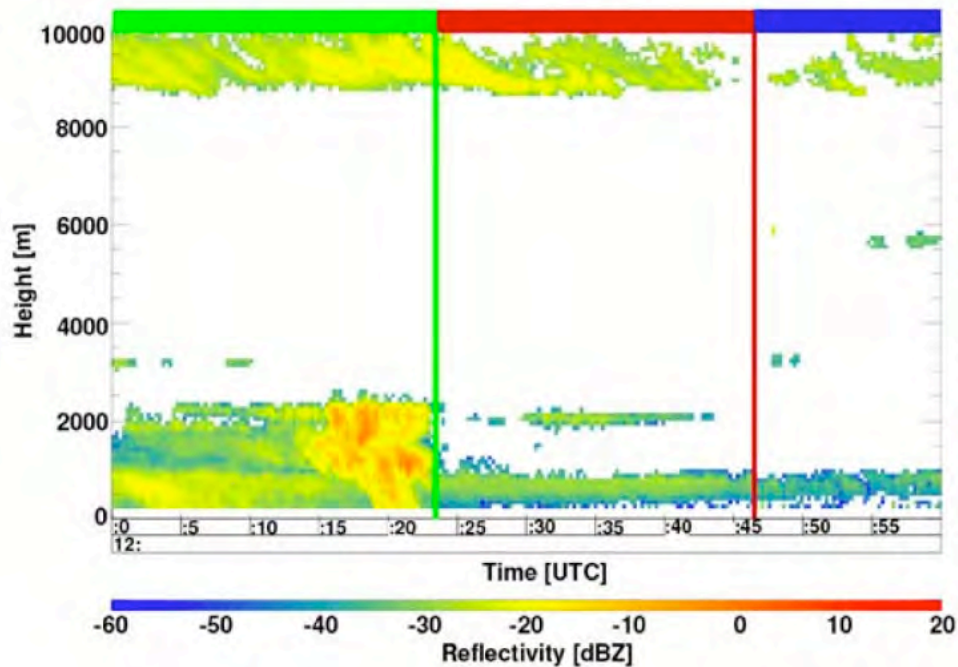
A complex cloud situation ( $1 < \alpha < 2$ ) with multi-layers, some broken; power-laws in  $\alpha-1$  will fit the data.



Min, Q.-L., L. C. Harrison, and E. E. Clothiaux, 2001: Joint statistics of photon path length and cloud optical depth: Case studies, *J. Geophys. Res.*, **106**, 7375-7385.

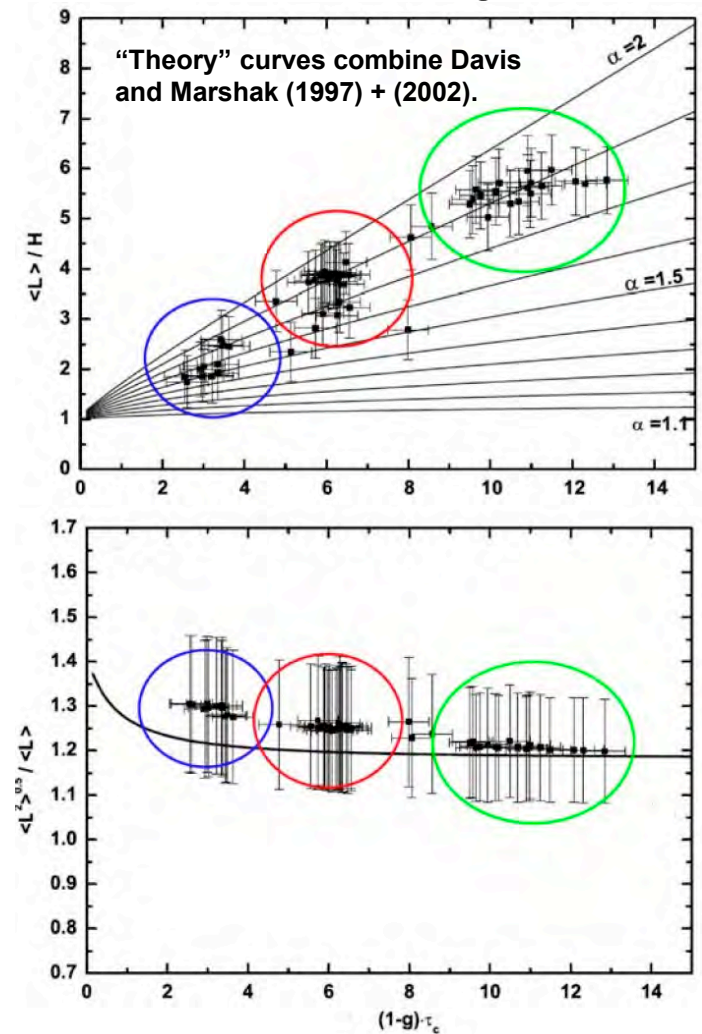
# From 1st to 2nd Moments ...

KNMI 35 GHz cloud radar

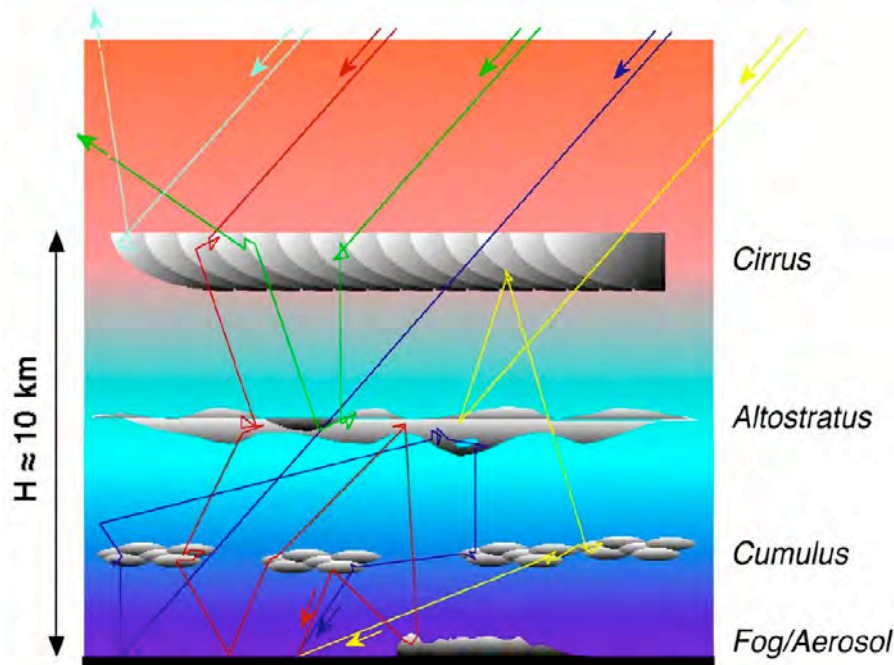


Scholl, T., K. Pfeilsticker, A. B. Davis, H. Klein Baltink, S. Crewell, U. Löhnert, C. Simmer, J. Meywerk, and M. Quante, 2006:  
 Path Length Distributions for Solar Photons Under Cloudy Skies: Comparison of Measured First and Second Moments with Predictions from Classical and Anomalous Diffusion Theories, *J. Geophys. Res.*, vol. 111, D12211-12226.

U. Heidelberg O<sub>2</sub> A-band spectrometer (2nd generation)



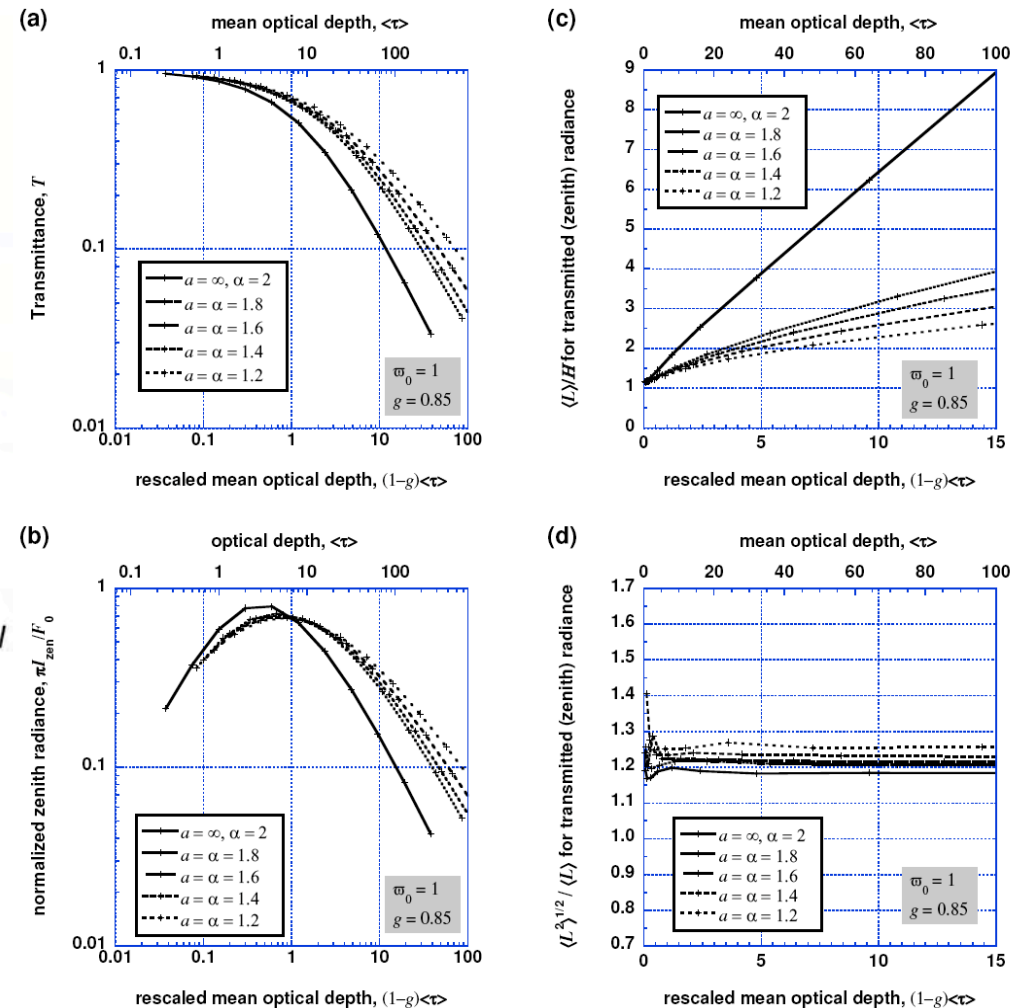
# Anomalous *Transport* Model



Solve 1D integral RTE with

$$T(s) = e^{-\sigma s} \mapsto$$

$$\langle T_a(s) \rangle = \frac{1}{\left(1 - \frac{\langle \sigma \rangle s}{a}\right)^a}, 1 < a \leq \infty$$



Davis, A. B., 2006: Effective Propagation Kernels in Structured Media with Broad Spatial Correlations, Illustration with Large-Scale Transport of Solar Photons Through Cloudy Atmospheres, in *Computational Methods in Transport – Granlibakken 2004* (Lecture Notes in Computational Science and Engineering), F. Graziani (Ed.), Springer-Verlag, New York (NY), pp. 85-140.

# 1. Build GCM SW transport schemes (1D-type RT models for averages over large domains)

- **Conceptualize unresolved variability**
  - Cloud fraction, cloud aspect ratio
  - Internal variability?
  - Spatial correlations, including layer-to-layer overlap
  - Need other parameters?
- **Make reasonable statistical assumptions**
  - Fractals and power-laws
  - Exponentials, lognormals, Gamma, etc.
- **Use judicious approximations in the RT**
- **Write code ... and *verify* it!**

## 2. Test GCM SW transport schemes with detailed 3D RT

- **Need 3D clouds from observations**
  - Could use CSRM output [OK for starts]
  - Better to use imagers and profilers
    - Ground-based: ARM sites
    - Space-based: A-train, especially CloudSat/Calipso
  - Best to use ARM Volume-imaging Array (AVA)
    - Need confidence in extracting optical properties
- **Need *verified* 3D RT codes**
  - full cloudy column capability
  - use to “assimilate” available data
- **Compare 3D (spatially averaged) and 1D-type RT model outputs for given inputs**



### **3. Validation of GCM grid-scale SW transport schemes, a.k.a. “closure” experiments**

- **Need new observations directly related to the domain-average HR profile predictions (i.e., output of new/improved SW transport codes)  
... at selected  $\lambda$ 's**
- **Focus on process (physics) of absorption by interstitial gases under complex cloudy conditions**
  - **use O<sub>2</sub> A-band:** known amount and known cross-section, but unknown path distribution
  - from below (ARM), within (UAVs), and above (OCO) ...
  - fully stress the 1D-type RT model: use scenarios with multiple and/or broken cloud layers

# ACT 4

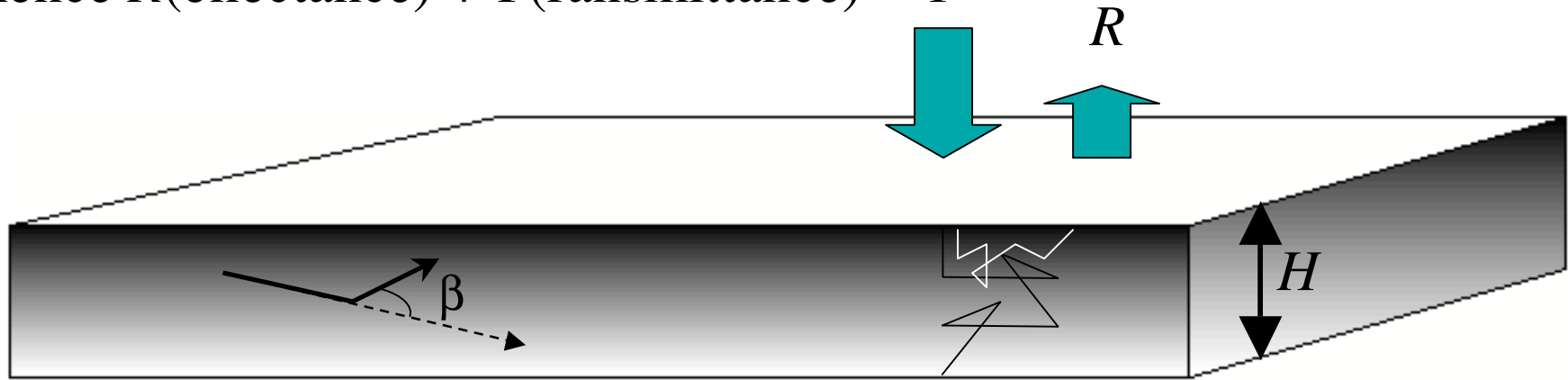
**Seeing clouds from both sides ...**





# Passive 3D cloud remote sensing in VNIR, 1

Take wavelength where there is no absorption,  
hence  $R(\text{reflectance}) + T(\text{transmittance}) = 1$

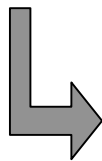


Mean Free Path  $\ell$



$$T = 1/[1+(1-g)\tau/2]$$

“Asymmetry” factor  $g = \langle \cos\beta \rangle \approx 0.85$



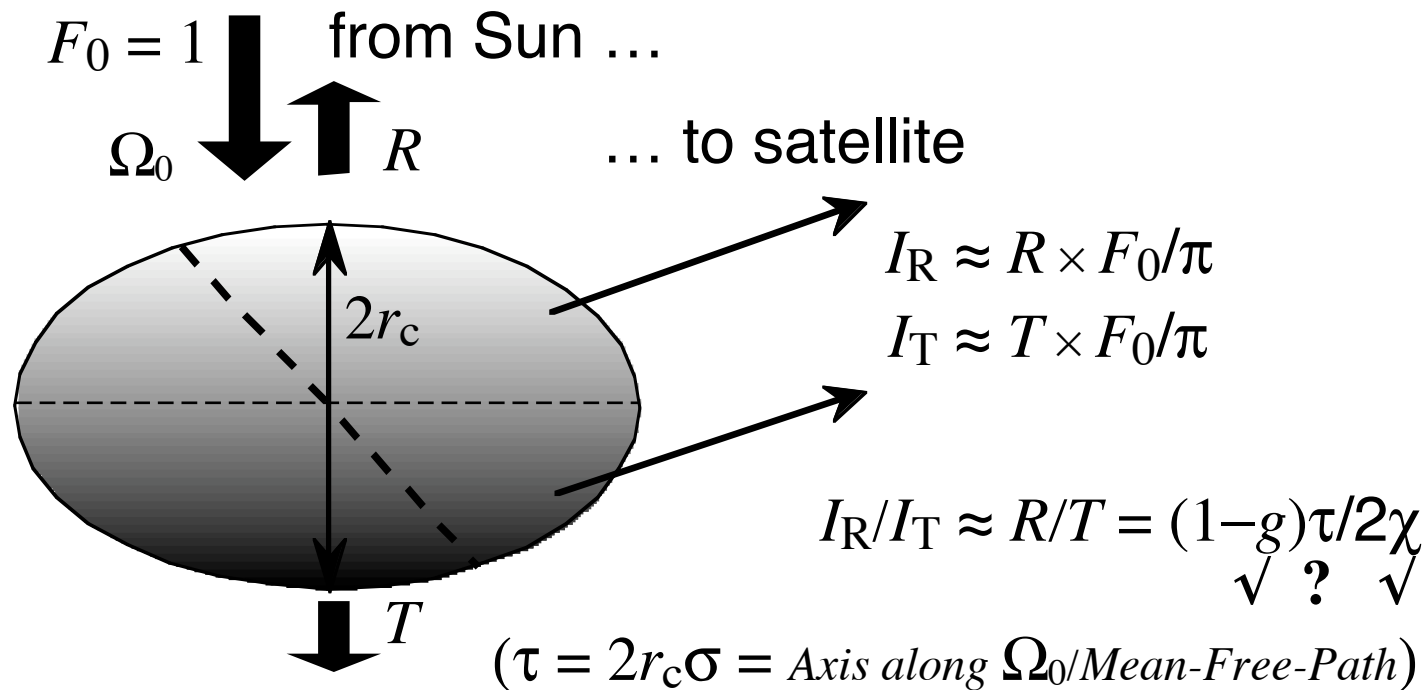
“Transport” MFP  $\ell_t = \ell/(1-g)$

“Optical” Depth  $\tau = H/\ell \gg 1$

“Rescaled” OD  $= H/\ell_t = (1-g)\tau$

Schuster, A., 1905: Radiation through a foggy atmosphere, *Astrophys. J.*, **21**, 1-22.

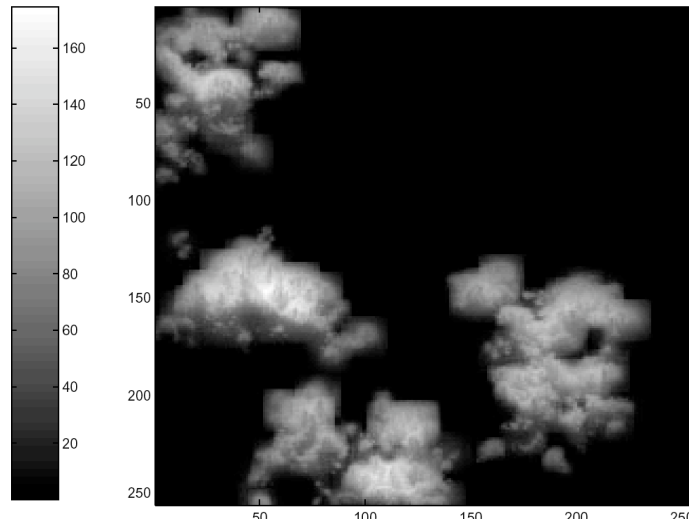
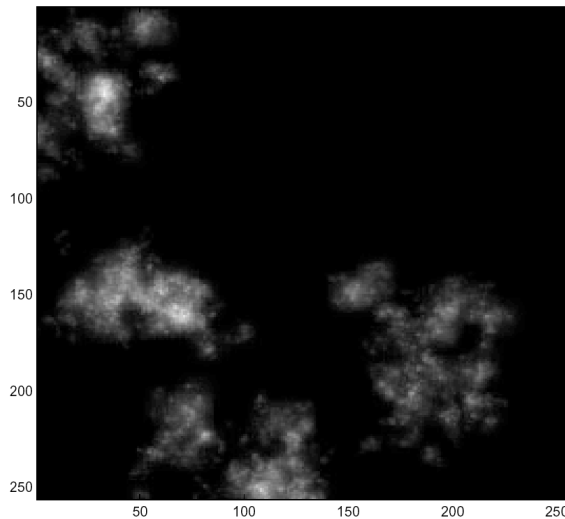
# Passive 3D cloud remote sensing in VNIR, 2



Davis, A. B., 2002: Cloud remote sensing with sideways-looks: Theory and first results using Multispectral Thermal Imager (MTI) data, in *S.P.I.E. Proceedings*, vol. 4725: “*Algorithms and Technologies for Multispectral, Hyperspectral, and Ultraspectral Imagery VIII*,” Eds. S. S. Shen and P. E. Lewis, S.P.I.E. Publications, Bellingham (Wa), pp. 397-405.

Unpublished generalizations with **Mathematica** and with **Igor Polonsky** (LANL PDF, now CSU).

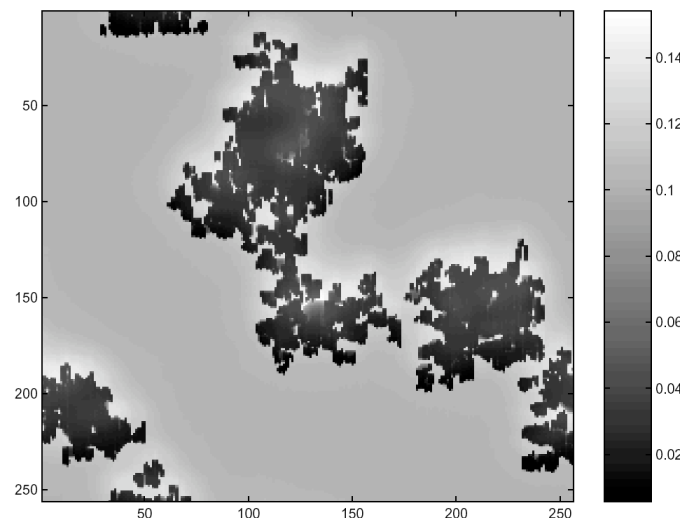
# Passive 3D cloud remote sensing in VNIR, 3



Will be used in validation of new  $\tau$  retrieval scheme (Polonsky et al., in preparation)

## Cloud Adjacency Effects, Simulated

- 3D clouds by wavelet-like “tdMAP” model (Benassi et al., 1999, 2004)
- 3D radiances by SHDOM code (Evans, 1998)



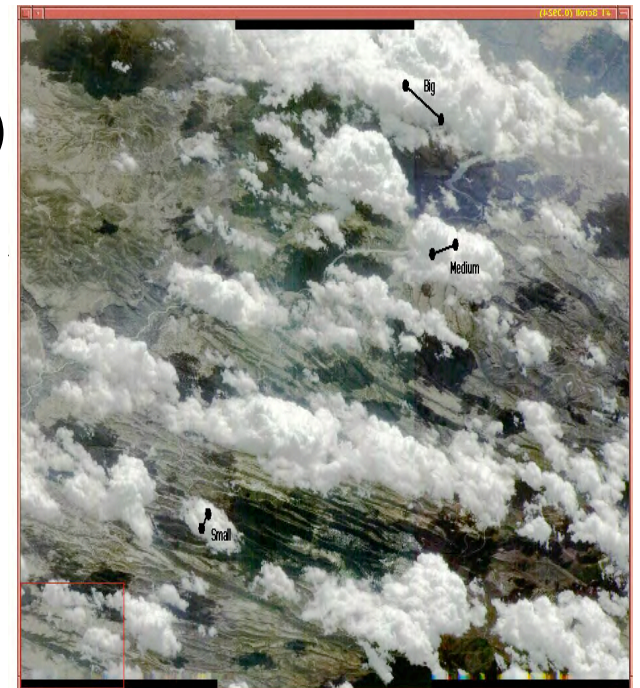
- Surface albedo assumed uniform at 0.1
- Appear to be as much as 50% more near clouds
- Depends on wavelength, hence confusion in spectral matching

## Passive 3D cloud remote sensing in VNIR, 4

- $\approx 100$  years after Schuster (1905),  $\approx 50$  after Chandrasekhar-Wick DOM in 1D, and  $\approx 25$  after 1D asymptotic theory ... a simple 3D RT solution for cloud remote sensing.

### Three applications areas:

- Hi-res imagery
  - Landsat, ASTER, MTI (DOE/NNSA)
  - ARM's Whole Sky Imagers?
- Cloud-capable atmospheric compensation for remote sensing of surface properties
- Effects of unresolved structure



# Epilogue



# Research Capabilities

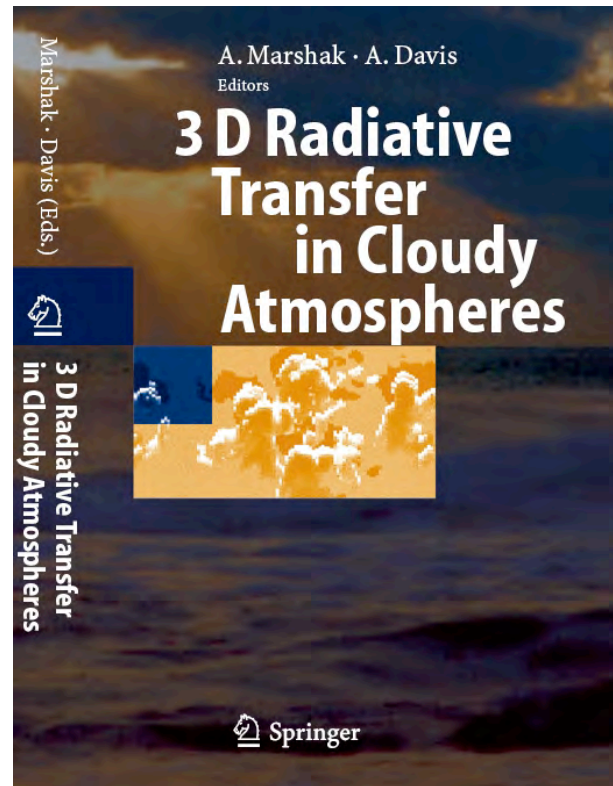
- **Atmospheric radiative transfer in or in the presence of clouds (i.e., dense structured media)**
  - Three-dimensional, for diagnostics & for energetics
  - Time-dependent, as needed
  - Theoretical (computational + analytical) and observational
- **Remote sensing techniques**
  - Passive and active, mostly in the optical (solar) spectrum
  - Physics-based *multi-pixel* exploitation methods
  - Instrument development
- **Concepts from nonlinear geophysics**
  - Multi-scale (wavelets, fractals, multifractals, etc.)
  - Stochastic modeling (cloud structure, etc.)
  - Model calibration, verification & (process of) validation

# Questions?





... for *some* answers:



(appeared in 2005)



# RT as a Linear Boltzmann Problem

- $I_\lambda(ct, \mathbf{x}, \Omega)$  is “radiance”  $\sim f(t, \mathbf{x}, \mathbf{p})$  for Boltzmann’s mesoscopic equation
  - $t$  is time and  $ct$  is total path (but we’ll consider stationary sources here)
  - $\mathbf{x}$  is position
  - $\Omega$  is direction of propagation (from  $\mathbf{p}/p$ )
  - $\lambda$  is wavelength (from  $h/p = ch/E$ ,  $E = h\nu$ )
- $\sigma(\mathbf{x})$  is the given “extinction” field = cross-section  $\times$  (fluctuating) density
  - $\sigma(\mathbf{x})$  is the collision probability per unit of length along a beam (we assume geometric optics)
  - $\sigma_s(\mathbf{x}) \leq \sigma(\mathbf{x})$  is the collision probability per unit of length for a scattering (rather than an absorption)
  - $\mathbf{x} \in M \subseteq \mathbb{R}^3$ , the optical medium
- $f_v(\mathbf{x})$ ,  $\mathbf{x} \in M$ , and  $f_s(\mathbf{x})$ ,  $\mathbf{x} \in \partial M$ , are the bulk and boundary emission rates.

$$I(\mathbf{x}, \Omega) = \int_0^{s_b(\mathbf{x}, -\Omega)} \sigma_s(\mathbf{x} - \Omega s') e^{-\int_0^{s'} \sigma(\mathbf{x} - \Omega s'') ds''} \left[ \int_{4\pi} p(\mathbf{x} - \Omega s', \Omega' \rightarrow \Omega) I(\mathbf{x} - \Omega s', \Omega') d\Omega' \right] ds' + Q(\mathbf{x}, \Omega)$$

with source term  $Q(\mathbf{x}, \Omega) = \int_0^{s_b(\mathbf{x}, -\Omega)} f_v(\mathbf{x} - \Omega s') e^{-\int_0^{s'} \sigma(\mathbf{x} - \Omega s'') ds''} ds' + f_s(\mathbf{x} - \Omega s_b(\mathbf{x}, -\Omega)) e^{-\int_0^{s_b(\mathbf{x}, -\Omega)} \sigma(\mathbf{x} - \Omega s') ds'}$

“propagation” kernel:  $T_0(s; \mathbf{x}, \Omega) = e^{-\int_0^s \sigma(\mathbf{x} - \Omega s') ds'}$   $= \exp[-\underbrace{\bar{\sigma}(s, \mathbf{x} + \Omega s / 2)}_{\text{random}} \times \underbrace{s}_{\text{fixed}}]$   $\dots \langle T_0(s) \rangle?$

# When $\sigma(x)$ is a function ...

$$\bar{\sigma}(r, x) = \frac{1}{r} \int_{x-r/2}^{x+r/2} \sigma(x') dx'$$

Empirical necessity to define a density (“1-point scale independence”):

$$\bar{\sigma}(r, x) \stackrel{d}{=} \bar{\sigma}(r', x)$$

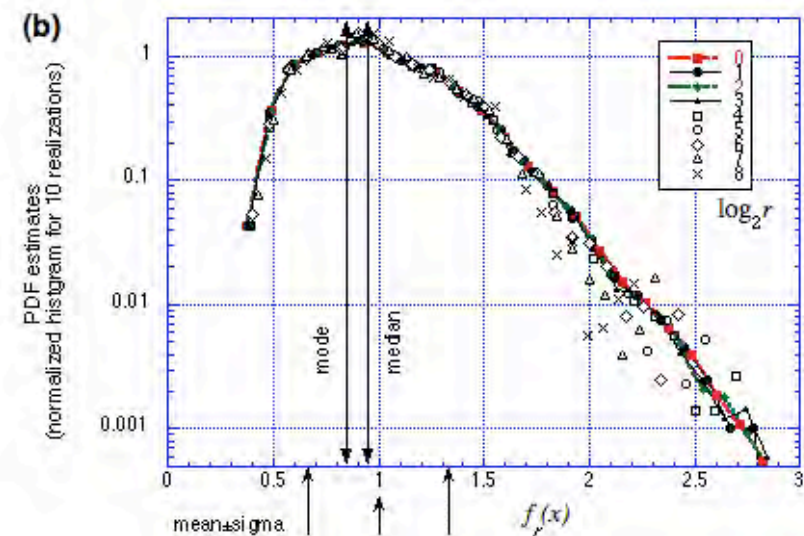
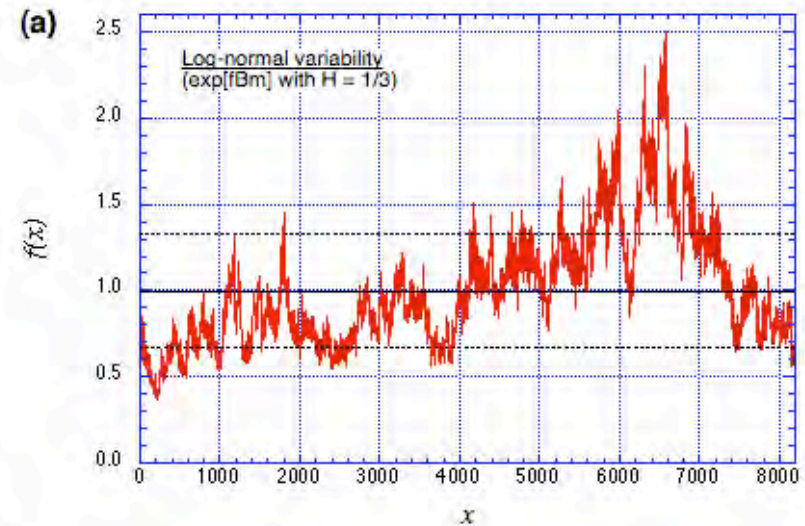
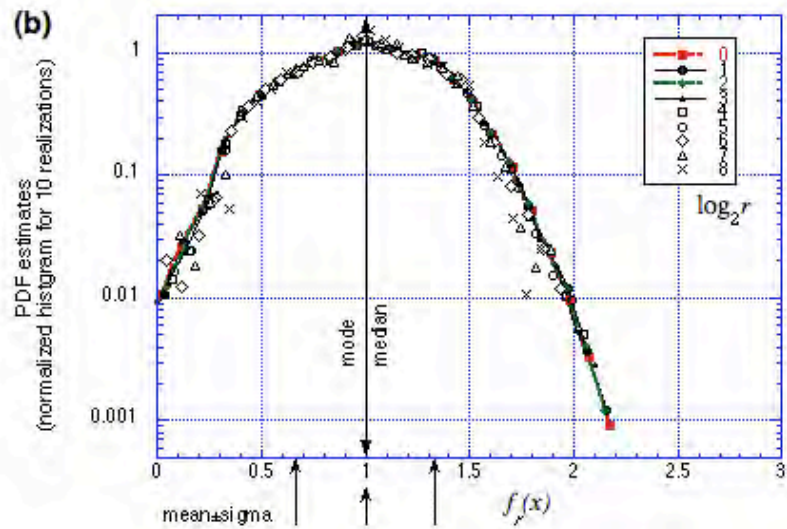
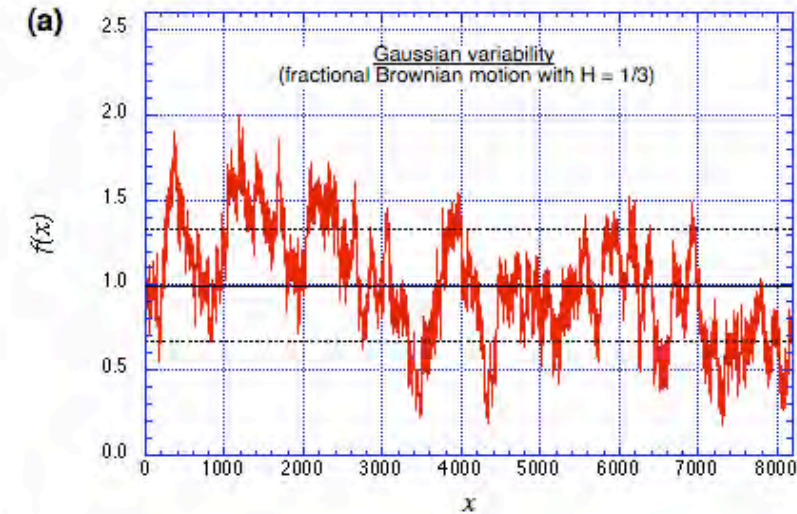
for  $r$  and  $r'$  in some large enough range.

$$\langle \sigma(x+r) - \sigma(x) \rangle \sim |r|^{h(x)}, 0 < h(x) \leq 1, \Rightarrow \lim_{r \rightarrow 0} \bar{\sigma}(r; x) = \sigma(x) \in \mathbb{R}^+, \forall x.$$

- I.  $\langle T_0(s) \rangle = \langle \text{Pr}\{\text{step} > s\} \rangle$  is exponential if and only if  $\sigma(x) \equiv \text{const.}$
- II. Mean-Free-Path (MFP) is minimal (and  $= 1/\sigma$ ) if  $\sigma(x) \equiv \text{const.}$
- III.  $\langle T_0(s) \rangle$  is sub-exponential, even if the real MFP is used (rather than  $1/\langle \sigma \rangle$ ).

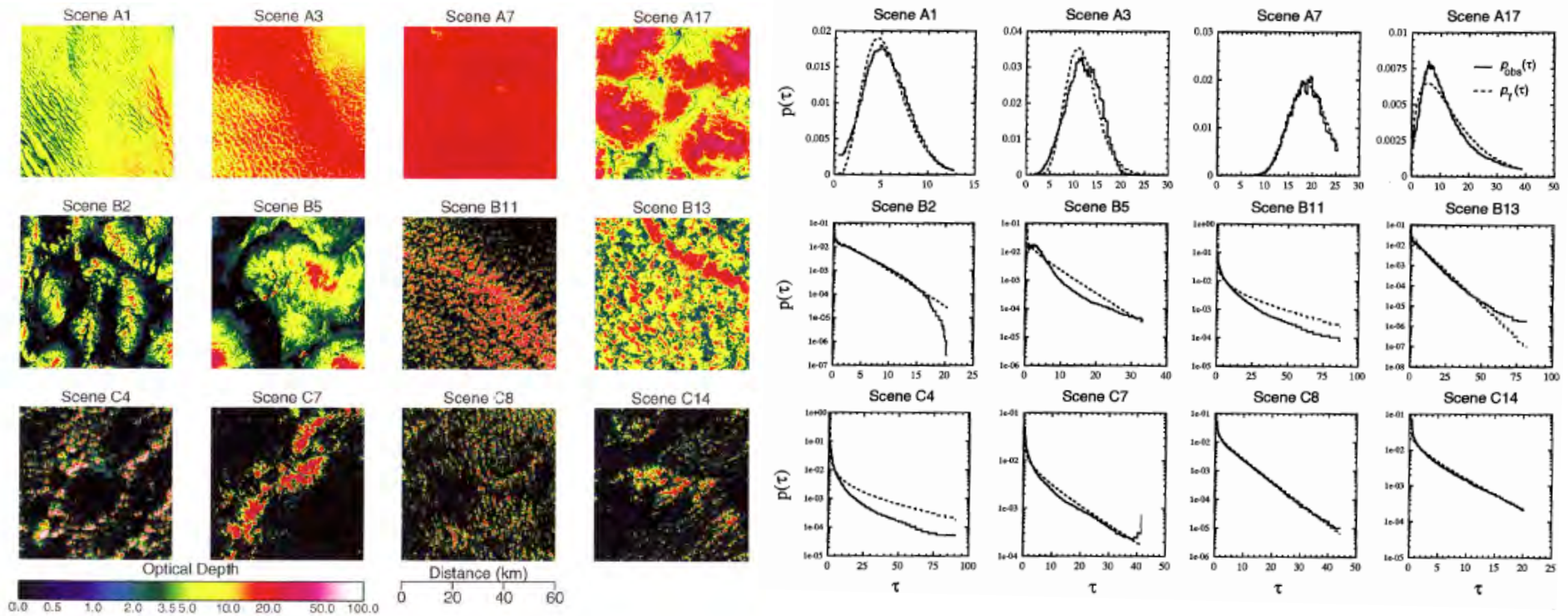
Davis, A. B., and A. Marshak, Photon propagation in heterogeneous optical media with spatial correlations: Enhanced mean-free-paths and wider-than-exponential free-path distributions, *J. Quant. Spectrosc. Rad. Transf.*, **84**, 3-34 (2004).

# Synthetic scale-invariant media that are turbulence-like



# Expectations for Earth's cloudy atmosphere, 1: *Barker et al.'s (1996) LandSat Data Analysis*

Gamma distributions capture many cloud optical depth scenarios.

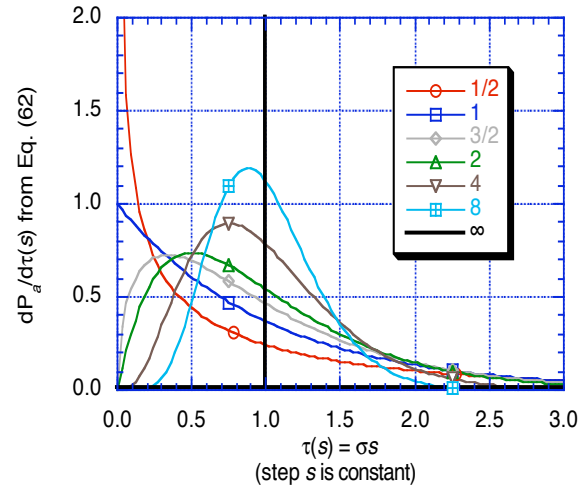


From: Barker, H. W., B. A. Wielicki, and L. Parker, 1996: A parameterization for computing grid-averaged solar fluxes for inhomogeneous marine boundary layer clouds - Part 2, Validation using satellite data, *J. Atmos. Sci.*, **53**, 2304-2316.



# Expectations for Earth's cloudy atmosphere, 2: *Effective transport kernels are power-law!*

(a) Gamma Probability Density Functions  
with  $\langle\sigma\rangle s = \langle\tau(s)\rangle = 1$ ,  $a = 1/\text{var}[\tau(s)]$



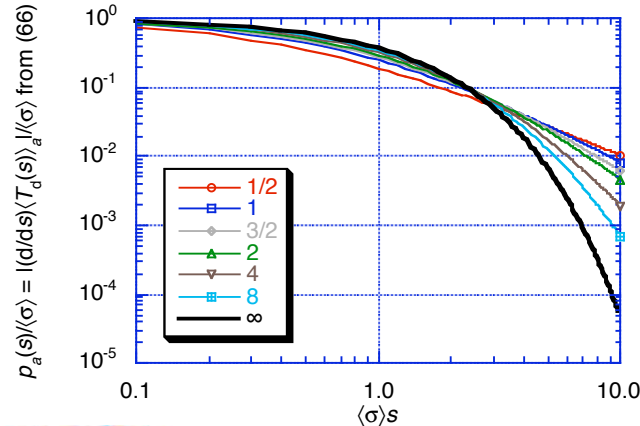
Assuming  $s = H$  (thickness) in previous slide:

$$p(\bar{\sigma}; \langle\bar{\sigma}\rangle, a) \approx \frac{1}{\Gamma(a)} \left( \frac{a}{\langle\bar{\sigma}\rangle} \right)^a \bar{\sigma}^{a-1} \exp[-a \frac{\bar{\sigma}}{\langle\bar{\sigma}\rangle}],$$

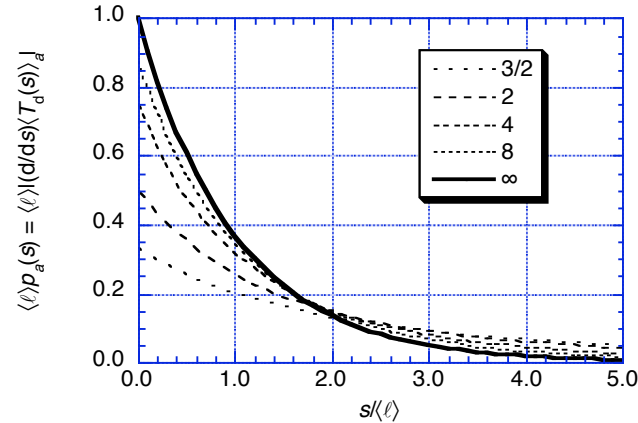
where  $a = \frac{1}{\langle\bar{\sigma}^2\rangle / \langle\bar{\sigma}\rangle^2 - 1}$ , yields mfp  $\ell = \left\langle \frac{1}{\bar{\sigma}} \right\rangle = \frac{a}{a-1} \frac{1}{\langle\bar{\sigma}\rangle}$

and  $\langle T_0(s) \rangle = \langle \exp[-\bar{\sigma}s] \rangle = \frac{1}{\left( 1 + \frac{\langle\bar{\sigma}\rangle s}{a} \right)^a} = \frac{1}{\left[ 1 + \frac{s/\ell}{(a-1)} \right]^a}.$

(a) Ensemble-averaged Free-Path Distributions (FPDs)  
for Gamma-distributed optical distances (fixed  $\langle\sigma\rangle$ )



(b) Ensemble-averaged Free-Path Distributions (FPDs)  
for Gamma-distributed optical distances (fixed  $\langle\ell\rangle$ )



# New 1D RT models that include the impact of 3D cloudiness

**Given:**

$$\begin{cases} I(z, \mu > 0) = \varpi_0 \int_0^z \left| \frac{d}{ds} \left\langle T_0 \left( s = \frac{z - z'}{\mu} \right) \right\rangle \right| \int_{4\pi} p(\mathbf{\Omega}' \rightarrow \mathbf{\Omega}) I(z', \mathbf{\Omega}') d\mathbf{\Omega}' \frac{dz'}{\mu} + Q(z, \mu) \\ I(z, \mu < 0) = \varpi_0 \int_z^H \left| \frac{d}{ds} \left\langle T_0 \left( s = \frac{z' - z}{|\mu|} \right) \right\rangle \right| \int_{4\pi} p(\mathbf{\Omega}' \rightarrow \mathbf{\Omega}) I(z', \mathbf{\Omega}') d\mathbf{\Omega}' \frac{dz'}{|\mu|} \end{cases}$$

where  $Q(z, \mu) = (\mu F_0 / \pi) \times \langle T_0(z/\mu) \rangle$  for  $\mu > 0$ , and 0 otherwise.

**Find:**

$$\begin{cases} T = \frac{2\pi}{F_0} \int_0^1 \mu I(z = H, \mu) d\mu \\ R = \frac{2\pi}{F_0} \int_0^1 |\mu| I(z = 0, -|\mu|) d|\mu| \end{cases}$$

$$I_{\text{zen}} = \pi I(z=H, \mu=+1) / F_0$$

$$I_{\text{nad}} = \pi I(z=0, \mu=-1) / F_0$$